# Meshing strategies for the alleviation of mesh-induced effects in cohesive element models

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#### Abstract

One of the main approaches for modeling fracture and crack propagation in solid materials is adaptive insertion of cohesive elements, in which line-like (2D) or surface-like (3D) elements are inserted into the finite element mesh to model the nucleation and propagation of failure surfaces. In this approach, however, cracks are forced to propagate along element boundaries, following paths that in general require more energy per unit crack extension (greater driving forces) than those followed in the original continuum, which in turn leads to erroneous solutions. In this work we illustrate how the introduction of a discretization produces two undesired effects, which we term mesh-induced anisotropy and meshinduced toughness. Subsequently, we analyze those effects through polar plots of the path deviation ratio (a measure of the ability of a mesh to represent straight lines) for commonly adopted meshes. Finally, we propose to reduce those effects through K-means meshes and through a new type of mesh, which we term conjugate-directions mesh. The behavior of all meshes under consideration as the mesh size is reduced is analyzed through a numerical study of convergence.

### 1 Introduction

The classical cohesive theory of fracture finds its origins in the pioneering works by Dugdale, Barenblatt and Rice [1, 2, 3]. In their work, fracture is regarded as a progressive phenomenon in which crack formation takes place across a cohesive zone ahead of the crack tip and is resisted by cohesive tractions. Cohesive zone models are widely adopted by scientists and engineers perhaps due to their straightforward implementation within the traditional finite element formulation. Some of the mainstream technologies proposed to introduce the cohesive

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theory of fracture into finite element analysis are the eXtended Finite Element Method (X-FEM) and cohesive elements.

Sukumar et al. [4] first utilized the X-FEM for modeling 3D crack growth by adding a discontinuous function and the asymptotic crack tip field to the finite elements. Subsequently, the method was extended to account for cohesive cracks [5]. It is worth noting that while the X-FEM approach can potentially deal with arbitrary crack paths, it becomes increasingly complicated for problems involving pervasive fracture and fragmentation.

On the other hand, the cohesive element approach consists on the insertion of cohesive finite elements along the edges or faces of the 2D or 3D mesh correspondingly [6, 7, 8, 9]. Even though this approach is well suited for problems involving pre-defined crack directions, a number of known issues affect the its accuracy when dealing with simulations including arbitrary crack paths, namely, (i) problems with the propagation of elastic stress waves (artificial compliance), (ii) spurious crack tip speed effects (lift-off), and (iii) mesh dependent effects (c.f. [10] for a comprehensive review). Despite these well known limitations, the robustness of the method makes it one of the most common approaches for pervasive fracture and fragmentation analysis.

Some of the limitations present in early approaches to cohesive element models were addressed by subsequent research efforts. For example, artificial compliance and lift-off effects can be avoided by using an initially rigid cohesive law [8] or, more elegantly, a discontinuous Galerkin formulation with an activation criterion for cohesive elements [11, 12]. However, the problem of mesh dependency is still an active area of research.

Mesh-dependent effects are direct consequence of cracks being able to propagate only across boundaries between bulk finite elements. That is, the topology of the mesh forces cracks to follow paths that in general require more energy per unit crack extension (greater driving forces) than those followed in the original continuum. In this work, we first focus on the effects that common topologies have on two main mesh-dependent effects, namely: mesh-induced toughness and mesh-induced anisotropy. We then illustrate how to decrease mesh-induced anisotropy through K-means meshes. Finally, we introduce a new type of mesh, termed conjugate-directions mesh, which greatly alleviate both effects.

The reminder of the paper is organized as follows: in section 2 we introduce a few preliminary concepts and discuss the mesh dependent effects; in section 3, we define K-means and conjugate-directions meshes and we study their mesh-dependent behavior, including a numerical convergence analysis; finally, in section 4, we summarize the main conclusions of the article.

#### 2 Mesh-induced effects

#### 2.1 Basic definitions

Let us consider a 2-dimensional brittle fracture problem in which an arbitrary crack of length  $L_c$  develops, as shown in Figure 1a. The fracture energy  $E_f$ 

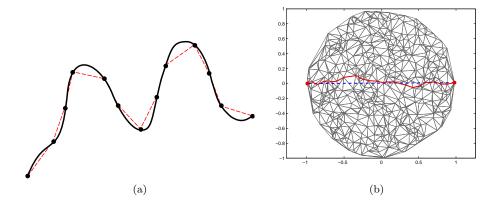


Figure 1: (a) Curve representing an arbitrary crack (solid line) and the corresponding rectification (dashed line); (b) shortest euclidean distance and shortest path over the mesh between two random points.

required to generate such a crack is proportional to the crack length  $L_c$ , that is,

$$E_f = L_c G_c \tag{1}$$

where  $G_c$  is the energy release rate of the material. This implies that, even assuming a correct model for the energy release rate of the material, the accuracy on any prediction of the fracture energy directly depends on a correct computation of the crack path length.

In the most general case, the crack path length could be computed by rectification of the irregular curve, i.e., the crack path could be approximated by connecting a finite number of points on the curve using line segments to create a polygonal path (see Figure 1a). The approximated crack length is computed as the sum of the length of the corresponding segments. The actual length could be then obtained by computing the limit of the approximated length as the segment's length tend to zero.

When a mesh is introduced to represent the continuum fracture problem within the cohesive element formulation, a constraint is introduced into the problem due to the inability of the mesh to represent, through its edges (2D) or faces (3D), the shape of an arbitrary crack. By following the reasoning in the preceding paragraph, a necessary condition for the convergence of the cohesive element approach is that the underlying mesh should be able to represent exactly a straight line (2D) or plane (3D) oriented in an arbitrary direction as the mesh is refined [13]. That is, instead of focusing our attention to the problem of representing arbitrary crack shapes, we could concentrate on the ability of the mesh to represent a straight line (2D) or a plane (3D), and then think of it as a necessary condition to be able to rectificate an arbitrary crack shape.

In 2-dimensional problems, the ability of a mesh to represent a straight

segment is characterized by the path deviation ratio  $\eta$ , defined as the ratio between the shortest path on the mesh edges connecting two nodes, and the Euclidian distance between them. The path deviation ratio can be interpreted as the maximum relative error in representing a straight line by means of the edges of a finite element mesh. Figure 1b illustrates the relevant quantities to be defined for the computation of the path deviation ratio on a given mesh. This particular mesh was generated by applying Delaunay's triangulation to a set of randomly distributed nodes. The dashed line represents the shortest path over the euclidean manifold between the two points marked in bold. The length of the dashed line is the euclidean distance between those points,  $L_e$ . On the same figure, the solid line represents the shortest path between the same pair of points over the edges of the mesh as computed by means of Dijkstra's algorithm [14]. The sum of the lengths of the segments composing the shortest path defines the length of shortest path on the graph between those points,  $L_g$ . According to the previous definitions, the path deviation can be computed as

$$\eta = \frac{L_g}{L_e} \tag{2}$$

Other associated measures are the absolute error in representing a straight line, given by  $\eta - 1$ , or the percentage error given by  $100 \times (\eta - 1)$ .

Finally, in order to compare path deviation ratios between different kinds of mesh and to compare meshes of different sizes, we introduce the non-dimensional mesh size, defined as

$$\lambda = \frac{\frac{1}{N} \sum_{i=1}^{N} h_i}{L_e} \tag{3}$$

where N is the total number of edges of the given mesh and  $h_i$  the length of the  $i^{th}$  edge. That is, the non-dimensional mesh size can be interpreted as the inverse of how many *average* segments are needed to represent a segment of length  $L_e$ .

#### 2.2 Mesh-induced anisotropy and mesh-induced toughness

By definition, the path deviation ratio is greater than one, i.e.,  $\eta \geq 1$ . Consequently, an arbitrary crack can only be approximated by a discrete curve of equal or larger length. For purposes of illustration, let us consider the 4k mesh depicted in Figure 2a. The diamond symbols on Figure 3a show, for such a mesh with  $\lambda \approx 1/200$ , how the absolute error changes as function of a chosen direction. As observed in the figure, the absolute error  $(\eta-1)=0$  for the horizontal, vertical, and  $\pm 45^{\circ}$  directions. However, it reaches a peak value near 0.08 for intermediate directions. In energetic terms, directions for which the absolute error (or path deviation ratio) is lower provide favorable paths for crack propagation. Clearly, cracks will tend to align along those lower-energy directions leading to inaccurate numeric predictions. We term this effect mesh-induced anisotropy.

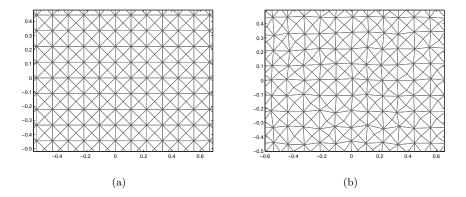
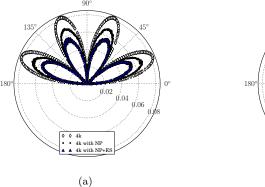


Figure 2: (a) typical 4k mesh, and (b) 4k mesh with nodal perturbation.

In general, when looking at all possible directions at a given point on the mesh, the expected value of the path deviation ratio is strictly larger than one, i.e.,  $\langle \eta \rangle > 1$ . This, in turn, implies that the expected value of the total energy released by the discrete crack  $\langle E_f \rangle_g$  is larger than the expected value of the energy  $E_f$  released by a real crack on the continuum. That is, the introduction of a discrete mesh necessarily leads to a larger effective toughness for the discrete cohesive model. We call this effect mesh-induced toughness.

A necessary condition to avoid the undesired mesh-induced anisotropy is for the path deviation ratio to be independent on the chosen direction. Additionally, a necessary condition to avoid the undesired mesh-induced toughness, is for the path deviation ratio to tend to 1 for all possible directions in the mesh as the non-dimensional mesh size tend to zero. The satisfaction of these two conditions is usually referred to as isoperimetric property. The work of Radin and Sadun [15] shows that the pinwheel tiling of the plane has the isoperimetric property. Based on this result, Papoulia et al. [13] observed that crack paths obtained from pinwheel meshes are more stable as meshes are refined when compared to other types of meshes. It is worth noting, however, that pinwheel meshes are hard to generate and there is no known extension to the 3-dimensional case.

Recently, Paulino et al. [16] introduced two mesh operators that decrease the undesired mesh dependent effects observed in 4k meshes. First, they introduced a nodal perturbation operator (NP operator) that randomly moves the nodes of 4k meshes by an amount proportional to the length of the shortest edge in the corresponding 4k cell. For example, Figure 2b shows a typical 4k mesh affected by the nodal perturbation operator. Subsequently, they introduced a topological edge swap operator (ES operator) that allows the originally vertical edges of the 4k mesh to switch to a horizontal position and vice versa. The idea behind these operators is to break the symmetry of the mesh at the geometrical level (NP operator) and to relax its structure at the topological level (ES operator).



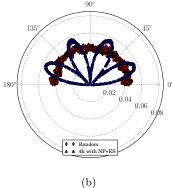


Figure 3: Error  $(\eta - 1)$  vs mesh direction for (a) 4k mesh, 4k mesh with nodal perturbation, and 4k mesh with nodal perturbation and edge swap, and (b) 4k mesh with nodal perturbation and edge swap, and random mesh.

As noted in their work, both operators produce an increment in the value of the path deviation ratio in the directions that showed no error for the 4k mesh, while at the same time they reduce it for the intermediate directions. They also show that the reduction of the path deviation ratio in the intermediate directions is larger than the increment produced in the the horizontal, vertical, and  $\pm 45^{\circ}$  directions. The combination of these effects results in an effective reduction of the mean value of the path deviation ratio  $\langle \eta \rangle$ . It is worth noting, however, that even though the mesh induced-toughness is reduced in this way (that is, the mean value of the path deviation ratio is reduced), the meshes under consideration still exhibit high anisotropy, only evident when considering the full polar plot of the absolute error (or path deviation ratio). We illustrate this effect by analyzing a 4k mesh with nodal perturbation factor  $NP_f = 0.3$  (4k with NP in Figure 3a), and a 4k mesh perturbed by the same amount with the extra possibility of swapping edges (4k with NP+ES in Figure 3a). As Figure 3a shows, for meshes with non-dimensional mesh size  $\lambda \approx 1/200$ , the maximum attained value of the path deviation ratio is decreased by the application of the NP operator. Furthermore, further application of the ES operator results in even lower values of the path deviation ratio. However, for all three cases the path deviation ratio highly depends on the direction under consideration. That is, both the NP and ES operators seem to have little effect on the mesh-induced anisotropy.

To conclude this section, we observe an interesting feature that arises from analyzing the polar behavior of a random mesh, i.e., a mesh generated through the application of Delaunay's triangulation algorithm to a set of randomly placed points. Figure 3b depicts the absolute error dependence on the direction for a random mesh and a 4K mesh subject to the NP and ES operators. The adopted

value of the nodal perturbation factor is NP<sub>f</sub> = 0.3, and both meshes have comparable values of  $\lambda \approx 1/200$ . It is interesting to note that whereas the 4k mesh exhibits a lower mean value of the absolute error ratio ( $\langle \eta - 1 \rangle \approx 0.37$  vs  $\langle \eta - 1 \rangle \approx 0.43$ ), the random mesh shows a much better isotropic behavior (std( $\eta - 1$ )  $\approx 0.002$  vs std( $\eta - 1$ )  $\approx 0.016$ ). Needles to say, ideally we would like to have a mesh that reflects a compromise between these two behaviors, that is, a low absolute error that remains constant for all possible directions.

## 3 K-means and conjugate-directions meshes

### 3.1 Isotropy through K-means meshes

Based on the previous observations, we wish to generate random meshes in order to achieve isotropy in the sense of the path deviation ratio. At first, purely random meshes look attractive: even though they generate path deviation ratios with mean value slightly higher than those obtained by 4k meshes under the effects of NP and ES operators, they have the advantage of not exhibiting any undesired anisotropy. However, purely random meshes are very irregular and present highly distorted triangles rendering them unfit for finite element computations. To address this issue, we propose to smooth purely random meshes by subjecting them to a K-means clustering algorithm. The K-means clustering algorithm partitions  $spn \cdot n$  observations into n clusters in which each observation belongs to the cluster with the nearest mean. Then, a so called K-means mesh can be obtained as follows:

- $spn \cdot n$  uniformly random points are clustered into n nodes.
- The resulting nodes are then triangulated by means of Delaunay's algorithm.

For the purposes of this work, K-means clustering was implemented through Lloyd's algorithm. For more details on Lloyd's algorithm and the relationship between K-means clustering and centroidal Voronoi tesselations, see [17]. Figure 4 shows a typical K-means mesh obtained obtained by clustering  $128 \cdot 512 = 65536$  random points into 512 nodes (spn = 128, n = 512).

To better illustrate the effect of the K-means clustering algorithm on random meshes, we demonstrate the influence of the clustering parameter spn on both the spatial node distribution, and on the triangle quality distribution of the resulting meshes.

Figure 5 shows the effect of the clustering parameter spn on the spatial distribution of mesh nodes and the corresponding Fourier spectrum. The Figure is organized in a matrix layout where rows (from top to bottom) correspond to nodes clustered with spn values of 1, 8, 32 and 128. The left column shows the spatial distribution of nodes, the center column the corresponding 2D Fourier power spectrum, and the right column the radially averaged power spectrum.

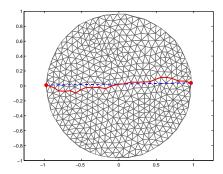


Figure 4: Typical K-means mesh obtained by clustering  $128 \cdot 512 = 65536$  random points into 512 nodes (spn = 128, n = 512).

The spn=1 case corresponds to a uniformly random distribution of nodes. The corresponding 2D-Fourier power spectrum exhibits an even frequency content typical of white noise. As the spn parameter increases, nodes tend to distribute more evenly in space, and the corresponding 2D-Fourier spectrum indicates that meshes remain isotropic (no radial variation in the spectrum) after the clustering algorithm. Due to this angular invariance, all spectral information can be exhamined through the radially averaged power spectrum. As it can be observed from the right column on Figure 5, the clustering algorithm filters the lower frequencies in the Fourier space, resulting in a typical blue noise power spectrum. That is, the lower the value of spn the closer the tendency to generate to a purely white noise mesh, where as spn increases the mesh approaches the behavior of a blue-noise mesh, which preserves isotropy and randomness while filtering low frequencies resulting in more even spacing between nodes.

Figure 6 shows the effect of the clustering parameter spn on the mesh quality distribution. The Figure is organized in a matrix layout where rows (from top to bottom) correspond to nodes clustered with spn values of 1, 8, 32 and 128. The left column shows the obtained mesh, and the right column the distribution of the mesh quality parameter q = 2r/R within the mesh, where r is the inradius and R is the circumradius of a triangle in the mesh. The chosen mesh quality parameter can take values between 0 and 1, 0 indicating a sliver (or zero-volume triangle) and 1 indicating an equilateral triangle. Clearly, as the value of spn increases the corresponding mesh quality distribution improves drastically. As a result, K-means meshes with higher value of the cluster parameter spn tend to contain a larger proportion of nearly equilateral triangles.

To conclude our study of K-means meshes, we analyze the effect of the clustering parameter spn on the metrics relevant to our original problem: the path deviation ratio. Figure 7a depicts the value of the mean value of the path deviation ratio as a function of spn for a mesh with non-dimensional mesh size  $\lambda \approx 1/400$ . As it can be observed, an increment on the value of spn

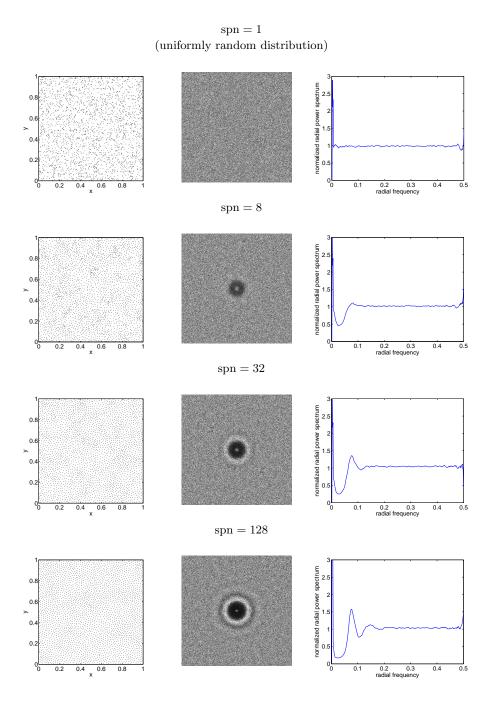


Figure 5: Influence of the smoothing parameter spn on mesh node distribution. Spatial node distribution (left column), 2D-Fourier power spectrum (center column), and radially averaged power spectrum (right column).

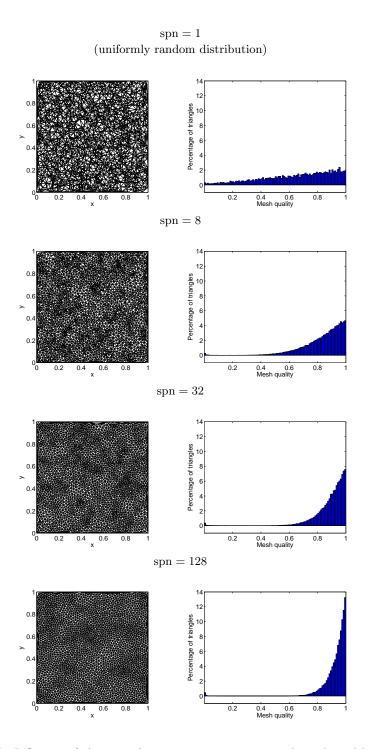


Figure 6: Influence of the smoothing parameter spn on mesh quality. Mesh (left column) and corresponding mesh quality distribution (right column).

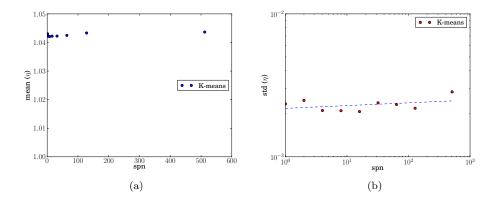


Figure 7: Influence of the smoothing parameter spn on (a) mean $(\eta)$ , and (b) std $(\eta)$  for a K-means mesh.

initially produces a slight reduction on the mean value of the path deviation ratio. After spn is further increased, the mean value of the path deviation ratio slightly increases, saturating at a value similar to that observed for the spn=1 case. Additionally, Figure 7b indicates that the standard deviation of the path deviation ratio is not much affected by the value of spn. That is, the smoothing process performed through the clustering algorithm improves the mesh quality while not affecting the behavior of path deviation ratio.

In summary, K-means meshes provide a good alternative to 4k meshes with NP and ES operators, as they provide slightly higher values for the mean value of the path deviation ratio while exhibiting a perfectly isotropic behavior.

### 3.2 Reducing mesh induced toughness through conjugatedirections meshes

The good isotropic behavior of the K-means meshes could be attributed to the ability of the mesh, at any node, to provide a set of random directions for the crack to continue its propagation towards a neighboring node. Similarly, the mean value of the path deviation ratio can be reduced by enriching the space of possible directions available at a point for a crack to propagate to the next point in the mesh, as done by the ES operator in 4k meshes. Towards this end, we introduce a new type of mesh, which we term conjugate-directions mesh. This new family of meshes preserves isotropy by using as starting point an existing K-means mesh, and enriches the space of possible directions by introducing topological modifications to it. Conjugate-directions meshes are obtained as follows:

• A K-means mesh is generated

#### A barycentric subdivision is applied to the initial K-means mesh

An algorithm for performing a barycentric subdivision to simplicial meshes of any dimension is detailed in [18], and explained in the subsequent paragraphs for completeness. Towards this end, we first describe (briefly) the concept of simplicial complex and subdivisions of a simplicial complex (triangulation).

A cell complex is a collection of objects, or cells, to which a precise dimension can be assigned. At its fundamental level, a 3-dimensional object can be represented by a collection of 0-cells (vertices), 1-cells (edges), 2-cells (faces) and 3-cells (volumes). In the subsequent paragraphs, p-dimensional cells are denoted as  $e_p$ . In addition,  $E_p(\mathcal{S})$  represents the collection of p-dimensional cells in a complex  $\mathcal{S}$ . The dimension dim  $\mathcal{S}$  of the cell complex  $\mathcal{S}$  is the largest dimension of any of its cells. Special classes of cell complexes are obtained when the cells are restricted to be of a certain type.

If all cells in a cell complex are *simplices*, the cell complex is said to be a *simplicial complex*. Simplicial complexes arise naturally as a result of the triangulation of solids. Specifically, a simplicial complex  $\mathcal{S}$  in  $\mathbb{R}^n$  is a collection of simplices in  $\mathbb{R}^n$  such that (i) every face of a simplex of  $\mathcal{S}$  is in  $\mathcal{S}$ , and (ii) the intersection of any two simplices of  $\mathcal{S}$  is a face of each of them.

We recall that the simplex  $\sigma$  spanned by a geometrically independent point set  $v_0, ..., v_p$  in  $\mathbb{R}^n$  is

$$\sigma = \left\{ x \in \mathbb{R}^n : \sum_{i=0}^p \lambda_i = 1, \sum_{i=0}^p \lambda_i v_i = x, \quad \lambda_i \ge 0 \quad \forall i = 0, 1, ..., p \right\}$$
 (4)

where the numbers  $\lambda_i$  are the barycentric coordinates of x with respect to  $v_0,...,v_p$ . The barycentric coordinates of a point are uniquely determined by it. The points  $v_0,...,v_p$  are the vertices of the simplex  $\sigma$  and p its dimension. Any simplex generated by a subset of  $v_0,...,v_p$  is a proper face of  $\sigma$  and the union of all proper faces is the boundary of the simplex. The notation  $e_\beta \prec e_\alpha$  signifies that  $e_\beta$  is a face of  $e_\alpha$ .

Let S be a cell complex. A subdivision complex of S is obtained by subdividing its cells into finer cells. More precisely, a complex  $S^*$  is said to be a subdivision of S if (i) every cell of  $S^*$  is contained in a cell of S, and (ii) every cell of S is the union of finitely many cells of  $S^*$ . These conditions particularly imply that the union of the cells of  $S^*$  equals the union of the cells of S and, hence,  $|S^*| = |S|$  as sets. The way in which a subdivision complex  $S^*$  is nested within the supercomplex S may be described by means of an inclusion map  $f: S^* \to S$ . This map assigns to every cell  $e_{\beta}$  of  $S^*$  the cell  $e_{\alpha}$  of S that contains it.

Barycentric subdivision generates a particular class of subdivision complexes. Thus, the barycentric subdivision of a simplicial complex is obtained by means of a uniform refinement. If  $\sigma_p = [v_0, ..., v_p]$  is a p-simplex of a simplicial complex  $\mathcal{S}$ , its barycenter is the point

$$\hat{\sigma}_p = \sum_{i=0}^p \frac{1}{p+1} v_i \tag{5}$$

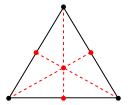


Figure 8: Barycentric subdivision of a triangle.

i. e., the barycenter is the point in the interior of  $\sigma_p$  that has equal barycentric coordinates. Given a simplicial complex S its barycentric subdivision  $S^*$  consists of all simplices  $[\hat{\sigma}_0, ..., \hat{\sigma}_p]$  such that  $\hat{\sigma}_0 \succ ... \succ \hat{\sigma}_p$ . The details of the baycentric subdivision algorithm are depicted in Algorithm 1 for completeness. As an example, Figure 8 shows the barycentric subdivision of a triangle. It is worth noting that the barycentric subdivision of a triangle will produce 6 new triangles whose aspect ratio is worse than that of the the original one.

```
Algorithm 1: Barycentric subdivision.
   Input: simplicial complex S, n = \dim S.
    Output: barycentric complex S^*.
    /* Initialize barycentric complex \mathcal{S}^* and inclusion map f
 1 for (p = 0, ..., n) do
       for (e_p \in E_p(S)) do

insert v in E_0(S^*);

insert (e_p, v) in f;
    /* Compute connectivity table CT(S)
                                                                                             */
 5 for (e_n \in E_n(\mathcal{S})) do
 6 insert [e_n] in CT(S);
 7 for (p = n, ..., 1) do
       for ([e_n, ..., e_p] \in CT(S)) do

for (e_{p-1} \prec e_p) do

append e_{p-1} to [e_n, ..., e_p];
    /* Compute connectivity table CT(S^*)
                                                                                             */
11 for ([e_n, ..., e_0] \in CT(\mathcal{S})) do
   insert [f(e_n), ..., f(e_0)] in CT(\mathcal{S}^*);
13 construct the simplicial complex S^* from CT(S^*);
14 return S^*
```

An interesting feature of the barycentric subdivision, and the reason why we chose it to generate conjugate-directions meshes, is that in the limiting case of equilateral triangles, the barycentric subdivision of a triangle provides an orthogonal direction for each direction provided by the edges of the original triangula-

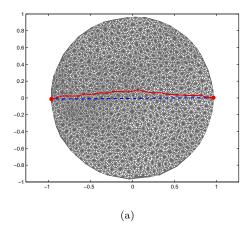


Figure 9: Typical conjugate-directions mesh obtained from a K-means mesh with spn = 128 and n = 512.

tion. This can be clearly observed in Figure 8, where there is a newly introduced dashed line perpendicular to each edge of the original triangle. In this way, for K-means meshes with high value of spn (for which we already illustrated their higher content of equilateral triangles), their barycentric subdivision enriches them with a set of directions that are almost orthogonal to the existing edges. Figure 9 shows a typical conjugate-directions mesh obtained from a K-means mesh with spn = 128 and n = 512.

Figure 10 shows the values of the absolute error vs direction for a 4k mesh with NP and ES, a K-means mesh, and a conjugate-directions mesh. All meshes are comparable in terms of their non-dimensional size  $\lambda \approx 1/200$ . The conjugate-directions mesh is not only isotropic, but the mean value of the absolute error is significantly lower ( $\approx 0.018$ ) than that observed for the other meshes ( $\approx 0.04$ ).

It is worth noting that, since conjugate-directions meshes are derived from K-means meshes, the spn parameter of the seed mesh will play a role in the performance of the newly generated one. Figure 11 shows the effect of the spn parameter on the path deviation ratio for a conjugate-directions mesh with  $\lambda \approx 1/400$ . As shown in 11a, an increment of spn produces a consistent reduction on the mean value of the path deviation ratio. For spn=512 the mean value of the path deviation ratio reaches a value  $\langle \eta \rangle \approx 1.015$ , approximately half (in error terms) of that observed for the barycentric mesh obtained from a random (spn=1) case. Additionally, Figure 11b indicates that the standard deviation of the path deviation ratio decreases as spn increases. This implies that, when a K-means mesh with high value of spn is adopted as a generator of a conjugate-directions mesh, the resulting triangulation produces a more compact distribution of the path deviation ratio.

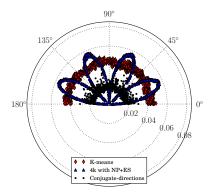


Figure 10: Error  $(\eta - 1)$  vs mesh direction for a 4k mesh with NP and ES, a K-means mesh, and a conjugate-directions mesh. All non-dimensional mesh sizes correspond to  $\lambda \approx 1/200$ .

In summary, conjugate-directions meshes provide an even better alternative to 4k meshes with NP and ES, as they provide significantly better values for the mean value of the path deviation ratio while being perfectly isotropic. It is worth noting that the original seed mesh must be of good quality since the barycentric subdivision stage in the generation of a conjugate-directions mesh produces a deterioration in the quality of the triangulation.

#### 3.3 Convergence of K-means and conjugate-direction meshes

To conclude this section, we perform a numerical convergence analysis of 4k, K-means and conjugate-direction meshes for the non-dimensional mesh size  $\lambda$  tending to zero. All K-means and conjugate-directions meshes considered in this numerical study of convergence were generated with a clustering parameter spn=128. The nodal perturbation factor adopted for all 4k meshes is NP<sub>f</sub> = 0.3.

Results show that convergence of K-means meshes, in the sense of the mean value of the path deviation ratio, is similar to that of 4k meshes with NP as the non-dimensional mesh size is reduced, see Figure 12a. Overall, 4k meshes with NP and ES exhibit lower mean values of the path deviation ratio within the mesh size range under consideration. This difference seems to decrease as the meshes are refined. Numerical evidence also shows that the mean value of the path deviation ratio tends to saturate around 1.04 for 4k with NP and K-means meshes, and around 1.036 for 4k mesh with NP and ES. That is, these three types of mesh have an *intrinsic roughness* that prevents them from representing straight lines with errors below 4% or 3.6% respectively no matter how refined they are.

On the other hand, the standard deviation of the path deviation ratio for

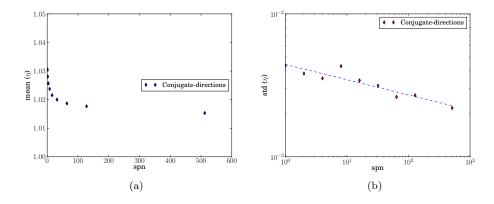


Figure 11: Influence of the smoothing parameter spn on (a)  $mean(\eta)$ , and (b)  $std(\eta)$  for a conjugate-directions mesh.

K-means meshes decreases at a fastest rate when compared to both types of 4k meshes, see Figure 12b. This indicates that even though both kinds of mesh exhibit similar mean values of the path deviation ratio for all mesh sizes, K-means meshes show less dispersion as the mesh is refined, while 4k meshes saturate very quickly. The saturation observed for both types of 4k mesh is an indicator that, in these cases, the mesh-induced anisotropy does not disappear as the meshes are refined.

In addition, the same numerical experiment shows no indication of saturation for conjugate-directions meshes in the studied range of non-dimensional mesh sizes. Furthermore, the mean value of the path deviation ratio is significantly smaller (on the order of 1/3 to 1/2 in absolute error terms) for conjugate-directions meshes when compared to the other meshes for all non-dimensional mesh sizes considered in this study. The standard deviation of the path deviation ratio seems to converge at a slower rate for this kind of mesh when compared to K-means meshes, but at a faster rate when compared to 4k meshes. It is worth noting, however, that it seems that the convergence rate on the standard deviation of the path deviation ratio tends to accelerate as the mesh is refined.

# 4 Summary and final remarks

We illustrated how the introduction of a discretization produces mesh-induced anisotropy and mesh-induced toughness when dealing with cohesive element modeling of arbitrary cracks. Subsequently, we analyzed those effects through the introduction of polar plots of the path deviation ratio, and showed that randomness plays a crucial role in the reduction of the former effect.

We observed that 4k meshes with NP and ES provide a better alternative to plain 4k meshes in the sense that they reduce the mesh-induced toughness.

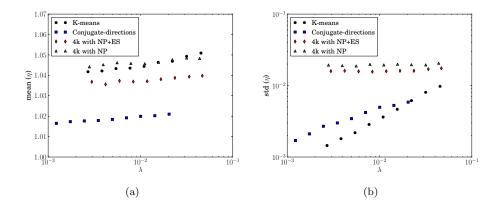


Figure 12: Numerical convergence analysis of 4k, K-means, and congugate-direction meshes. Mean value of path deviation ratio (a) and standard deviation of path deviation ratio (b) as a function of the non-dimensional mesh size.

However, when looking at the full polar plot of the path deviation ratio vs mesh direction, it becomes evident that mesh-induced anisotropy, while reduced, still persists. That is, regardless of the type of 4k mesh (with or without NP and/or ES), there is always an associated anisotropy to it. In addition, this effect seemingly does not disappear as the mesh is refined. This has profound implications on the interpretation of mesh refinement algorithms involving 4k meshes: at some point, there is no gain from mesh refinement no matter how much these meshes are refined.

On the other hand, K-means meshes were shown to be a good compromise between isotropy and good element quality. They are isotropic thus not providing preferred crack propagation directions for any mesh size. However, K-means meshes saturate, in the sense of the mean value of the path deviation ratio, as they are refined. This saturation results in the introduction of mesh-induced toughness even in the limiting case of non-dimensional mesh size tending to zero. A good feature of K-means meshes is that they do not seem to saturate in the sense of the standard deviation of the path deviation ratio, which implies that the error concentrates around its mean as the mesh is refined. This would allow to adopt K-means meshes in cohesive element modeling as long as the proper correction term is introduced in the toughness of the material.

Finally, we proposed a new type of mesh, termed conjugate-directions mesh, which significantly reduces the undesired mesh-dependent effects for meshes of practical size. Our approach combines the use of barycentric subdivision with K-means meshes to provide a greater distribution of directions for failure surfaces to propagate. Numerical evidence suggest that, due to the exhibited isotropy and considerably lower mesh-induced toughness effect, conjugate-direction meshes could be good candidates for cohesive element analysis of crack

propagation problems where crack paths are not known a priori. Our hope is that, by designing meshes to mitigate the effects of mesh topology, we can invoke adaptive insertion along pre-defined element boundaries to achieve arbitrary crack propagation. The great promise of this research effort is that arbitrary crack propagation could be achieved through mesh design (pre-processing).

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#### References

- [1] D. S. Dugdale. Yielding of steel sheets containing slits. *Journal of the Mechanics and Physics of Solids*, 8(2):100–104, 1960.
- [2] G. I. Barenblatt. The mathematical theory of equilibrium cracks in brittle fracture. Advances in Applied Mechanics, 7:55–129, 1962.
- [3] J. R. Rice. Mathematical analysis in the mechanics of fracture. Fracture: an advanced treatise, 2:191–311, 1968.
- [4] N. Sukumar, N. Moes, B. Moran, and T. Belytschko. Extended finite element method for three-dimensional crack modeling. *International Journal for Numerical Methods in Engineering*, 48:1549–1570, 2000.
- [5] N. Moes and T. Belytschko. Extended finite element method for cohesive crack growth. *Engineering Fracture Mechanics*, 69:813–833, 2002.
- [6] X. P. Xu and A. Needleman. Numerical simulations of fast crack growth in brittle solids. *Journal of the Mechanics and Physics of Solids*, 42:1397– 1397, 1994.
- [7] X. P. Xu and A. Needleman. Numerical simulations of dynamic crack growth along an interface. *International Journal of Fracture*, 74:289–324, 1995.
- [8] G. T. Camacho and M. Ortiz. Computational modelling of impact damage in brittle materials. *International Journal of Solids and Structures*, 33:2899–2938, 1996.
- [9] M. Ortiz and A. Pandolfi. Finite-deformation irreversible cohesive elements for three-dimensional crack-propagation analysis. *International Journal for Numerical Methods in Engineering*, 44(9):1267–1282, 1999.
- [10] A. Seagraves and R. Radovitzky. Advances in cohesive zone modeling of dynamic fracture. Springer Verlag, 2009.

- [11] L. Noels and R. Radovizky. An explicit discontinuous galerkin method for non-linear solid dynamics: Formulation, parallel implementation and scalability properties. *International Journal for Numerical Methods in En*gineering, 74:1393–1420, 2008.
- [12] R. Radovizky, A. Seagraves, M. R. Tupek, and L. Noels. A scalable 3d fracture and fragmentation algorithm based on a hybrid, discontinuous galerkin, cohesive element method. Computer Methods in Applied Mechanics and Engineering, 200:326–344, 2011.
- [13] K. D. Papoulia, S. A. Vavasis, and P. Ganguly. Spatial convergence of crack nucleation using a cohesive finite-element model on a pinwheel-based mesh. *International Journal for Numerical Methods in Engineering*, 67:1– 16, 2006.
- [14] E. W. Dijkstra. A note on two problems in connexion with graphs. *Numerische mathematik*, 1(1):269–271, 1959.
- [15] C. Radin and L. Sadun. The isoperimetric problem for pinwheel tilings. Communications in Mathematical Physics, 177(1):255–263, 1996.
- [16] G. H. Paulino, K. Park, W. Celes, and R. Espinha. Adaptive dynamic cohesive fracture simulation using nodal perturbation and edge-swap operators. International Journal for Numerical Methods in Engineering, 84(11):1303–1343, 2010.
- [17] Q. Du, V. Faber, and M. Gunzburger. Centroidal voronoi tessellations: Applications and algorithms. *SIAM Review*, 41:637–676, 1999.
- [18] J. J. Rimoli and M. Ortiz. A duality-based method for generating geometric representations of polycrystals. *International Journal for Numerical Methods in Engineering*, 86(9):1069–1081, 2011.